

TFY4215/FY1006

Delta function

Dirac's delta function can be loosely defined as

$$\delta(x) = \begin{cases} 0, & x \neq 0, \\ \infty, & x = 0. \end{cases} \quad (1)$$

in such a way that the area under the curve is constant

$$\int_{-\infty}^{\infty} \delta(x) dx = 1. \quad (2)$$

One can formally show that no such function exists. However, one can define a sequence of integrable functions $f_n(x)$, such that

$$\int_{-\infty}^{\infty} f_n(x) dx = 1, \quad (3)$$

for all n and such that $f_n(x)$ becomes increasingly concentrated around $x = 0$ as n increases. We can then define the integral of the δ -function by

$$\int_{-\infty}^{\infty} \delta(x) dx \equiv \lim_{n \rightarrow \infty} \int_{-\infty}^{\infty} f_n(x) dx. \quad (4)$$

A sequence of functions with these properties is called a δ -sequence. For example the sequence of Gaussians

$$f_n(x) = \frac{n}{\sqrt{\pi}} e^{-n^2 x^2}, \quad (5)$$

has these properties. This is illustrated in Fig. 1.

Another example is the sequence of functions $g_n(x)$ defined by

$$g_n(x) = \begin{cases} \frac{n}{2}, & |x| < \frac{1}{n}, \\ 0, & |x| \geq \frac{1}{n}. \end{cases} \quad (6)$$

For proofs, the sequence $g_n(x)$ is easier to use.¹

¹Note that all identities and relations must be independent of the particular sequence one is using. This can be shown rigorously.

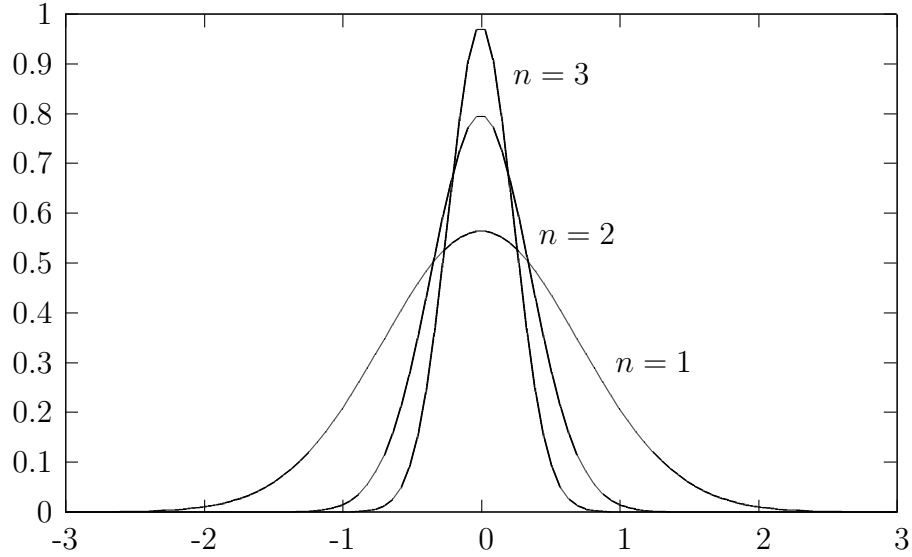


Figure 1: Sequence $f_n(x)$ of Gaussians which are increasingly concentrated around $x = 0$.

We are interested in calculating the integral

$$I \equiv \int_{-\infty}^{\infty} f(x) \delta(x) dx , \quad (7)$$

where $f(x)$ is some integrable function. Heuristically speaking, the δ -function picks up a contribution to the integral only when $x = 0$, and so our intuition tells us that the value of the integral should be $I = f(0)$. Let us confirm this expectation. In analogy with the definition Eq. (4), we define the integral I by

$$I \equiv \lim_{n \rightarrow \infty} \int_{-\infty}^{\infty} f(x) g_n(x) dx \quad (8)$$

Using the sequence (6), we obtain

$$\begin{aligned} \int_{-\infty}^{\infty} f(x) g_n(x) dx &= \frac{n}{2} \int_{-1/n}^{1/n} f(x) dx \\ &= f(\xi) , \end{aligned} \quad (9)$$

where $-\frac{1}{n} < \xi < \frac{1}{n}$. The last line follows from the mean-value theorem of calculus. In the limit $n \rightarrow \infty$, the value is $\xi = 0$ and so we obtain the important result

$$\boxed{\int_{-\infty}^{\infty} \delta(x) f(x) dx = f(0)} \quad (10)$$

Similarly, let us calculate the integral

$$I = \int_{-\infty}^{\infty} f(x) \delta(x - a) dx , \quad (11)$$

where a is a nonzero constant. Changing variable to $y = x - a$, we see that we must have $I = f(a)$:

$$\begin{aligned} I &= \int_{-\infty}^{\infty} f(y + a) \delta(y) dy \\ &= f(a) . \end{aligned} \quad (12)$$

We next consider the integral

$$I = \int_{-\infty}^{\infty} \delta(f(x)) dx , \quad (13)$$

If the function $f(x)$ is nonzero for all $x \in (-\infty, \infty)$, the integral vanishes. Assume next that the function has a single zero at $x = x^*$. Moreover assume that $f'(x^*) > 0$ so that $f(x)$ is invertible in a small interval around x^* . Introducing the variable $y = f(x)$, the integral can be written as

$$\begin{aligned} I &= \int_{x^* - \Delta}^{x^* + \Delta} \delta(f(x)) dx \\ I &= \int_{f(x^*) - f'(x^*)\Delta}^{f(x^*) + f'(x^*)\Delta} \delta(y) \frac{dx}{dy} dy \\ &= \int_{-f'(x^*)\Delta}^{f'(x^*)\Delta} \delta(y) [f^{-1}(y)]' dy \\ &= [f^{-1}(y)]' \Big|_{y=0} \\ &= \frac{1}{f'(x^*)} , \end{aligned} \quad (14)$$

where we in the last line have used $x = f^{-1}(y)$ and therefore $1 = [f^{-1}(y)]' f'(x)$. If $f'(x^*) < 0$, the integral changes sign, $I = -1/f'(x^*)$. The two cases can conveniently be summarized in the formula

$$I = \frac{1}{|f'(x^*)|} . \quad (15)$$

If the function $f(x)$ has more than one zero, one must simply sum over all the zeros x_i^* of $f(x)$ and we can write

$$\boxed{\delta(f(x)) = \sum_i \frac{1}{|f'(x_i^*)|} \delta(x - x_i^*)} \quad (16)$$

Example

Consider the following integral

$$I = \int_{-\infty}^{\infty} \delta(x^2 - a^2) dx . \quad (17)$$

We define $f(x) = x^2 - a^2$ and we then have $f'(x) = 2x$. The zeros of $f(x)$ are $x^* = \pm a$ and hence $|f'(x = x^*)| = 2|a|$. This implies that we can write

$$\begin{aligned} I &= \frac{1}{2|a|} \int_{-\infty}^{\infty} [\delta(x - a) + \delta(x + a)] dx \\ &= \frac{1}{|a|} . \end{aligned} \quad (18)$$

Note that the function $\delta(x^2)$ is too singular to be well defined. The problem is that the derivative of $f(x) = x^2$ vanishes at $x = 0$ which is the zero of the function itself and so Eq. (15) makes no sense. This also follows directly from Eq. (18).

Heaviside's step function

Let us define the function $\theta(x)$ by

$$\theta(x) = \int_{-\infty}^x \delta(x) dx . \quad (19)$$

Using the rules of delta-function calculus, we find ²

$$\theta(x) = \begin{cases} 1 , & x > 0 , \\ 0 , & x < 0 . \end{cases} \quad (20)$$

which is nothing but Heaviside's step function, see Fig. 2.

From the fundamental theorem of calculus, one obtains

$$\theta'(x) = \delta(x) . \quad (21)$$

Eq. (21) is a formal result as the derivative of Eq. (20). $\theta'(x)$ is not defined for $x = 0$.

²The defining Eq. (19) may not hold for $x = 0$. It depends on the limiting procedure that defines the integral of $\delta(x)$. For example using one of the delta sequences from this chapter, it is clear that $\theta(0) = 1/2$, which is also a common definition.

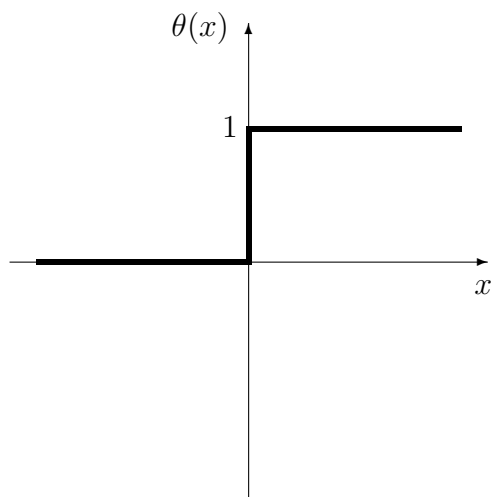


Figure 2: Heaviside step function.

Problems

- 1) Calculate the Fourier transform $f(p)$ of the delta function $\delta(x)$. Use this result to find an integral representation of $\delta(x)$.
- 2) Prove the *derivative sifting formula*:

$$\int_{-\infty}^{\infty} \delta'(x) f(x) dx = -f'(0) , \quad (22)$$

where $\delta'(x)$ is the derivative of the δ function and $f(x)$ is any differentiable function. Hint: use integration by parts and a *differentiable* δ -sequence, e.g. the Gaussian in Eq. (5).